

① Qubit density operator :-

Recall,  $\left\{ \frac{\mathbb{I}}{2}, \frac{\sigma_x}{2}, \frac{\sigma_y}{2}, \frac{\sigma_z}{2} \right\}$  form an ON basis for  $M_2(\mathbb{C})$ .

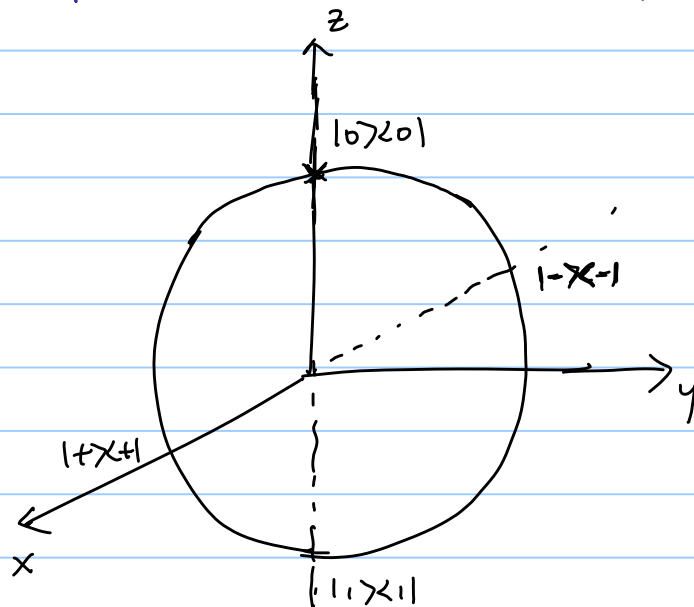
$$\therefore \rho = \frac{1}{2} (r_0 \mathbb{I} + r_x \sigma_x + r_y \sigma_y + r_z \sigma_z)$$

$$\text{Tr}(\rho) = 1 \Rightarrow r_0 = 1$$

$$\rho^\dagger = \rho \Rightarrow (r_x, r_y, r_z) \equiv \vec{r} \in \mathbb{R}^3.$$

$$\text{Tr}(\rho^2) \leq 1 \Rightarrow |\vec{r}| \leq 1$$

Unit sphere in 3-d  $\Rightarrow$  Bloch Sphere!!



$$\text{Origin: } \rho = \frac{1}{2} |0\rangle\langle 0| + \frac{1}{2} |1\rangle\langle 1|$$

$$= \frac{1}{2} |+\rangle\langle +| + \frac{1}{2} |-\rangle\langle -|$$

$\Rightarrow$  Maximally mixed state!

## ② Composite Systems & Entanglement

$$|\psi_A\rangle \in H_A, |\psi_B\rangle \in H_B.$$

$H_A, H_B$   
are called  
subsystems  
of  $H_{AB}$

$$|\Psi_{AB}\rangle \in H_A \otimes H_B.$$

↓  
Kronecker product

$$A \otimes B = \begin{pmatrix} a_{11}B & a_{12}B & \dots & a_{1n}B \\ \vdots & \vdots & \ddots & \vdots \\ a_{m1}B & a_{m2}B & \dots & a_{mn}B \end{pmatrix}$$

eg. Basis for  $\mathbb{C}^2 \otimes \mathbb{C}^2$  :-

$$\{|0\rangle, |1\rangle\}_A \otimes \{|0\rangle, |1\rangle\}_B.$$

$$\equiv \{|0\rangle \otimes |0\rangle, |0\rangle \otimes |1\rangle, |1\rangle \otimes |0\rangle, |1\rangle \otimes |1\rangle\}$$

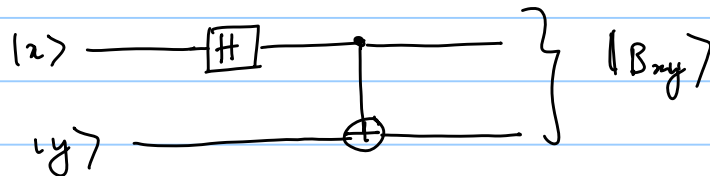
\* Rules:

$$* \langle \psi_A \otimes \psi_B, e_A \otimes f_B \rangle = \langle \psi_A, e_A \rangle \langle \psi_B, f_B \rangle$$

$$* (A \otimes B) (|\psi\rangle \otimes |\omega\rangle) = A|\psi\rangle \otimes B|\omega\rangle$$

\* Bell states:-

Consider this circuit:-



$$|\beta_{00}\rangle = \frac{|00\rangle + |11\rangle}{\sqrt{2}}, \quad |\beta_{01}\rangle = \frac{|01\rangle + |10\rangle}{\sqrt{2}}$$

$$|\beta_{10}\rangle = \frac{|00\rangle - |11\rangle}{\sqrt{2}}, \quad |\beta_{11}\rangle = \frac{|01\rangle - |10\rangle}{\sqrt{2}}$$

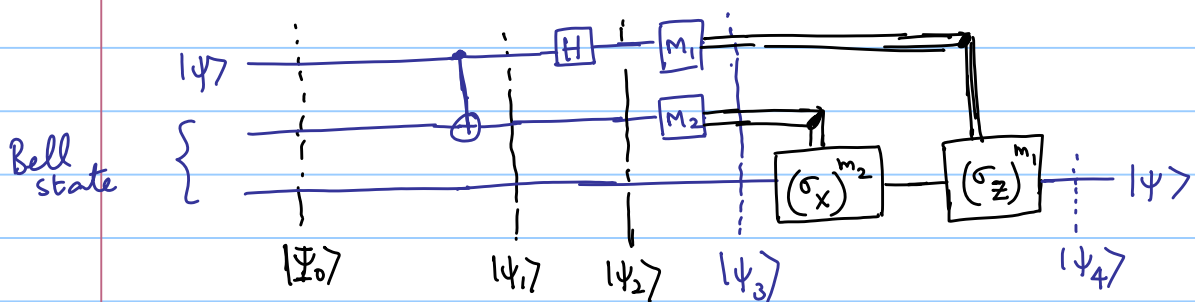
$\{|\beta_{00}\rangle, |\beta_{01}\rangle, |\beta_{10}\rangle, |\beta_{11}\rangle\}$  forms an ON basis for  $\mathbb{C}^2 \otimes \mathbb{C}^2$ .

\* There do not exist states  $|\alpha\rangle, |\beta\rangle \in \mathbb{C}^2$  such that

$$|\beta_{xy}\rangle = |\alpha\rangle \otimes |\beta\rangle$$

$\Rightarrow$  These are not product states; entangled!

(2a) Teleportation: - Entangled states are an important resource for quantum information and computing.



$|\psi\rangle = \alpha|0\rangle + \beta|1\rangle$  ("unknown" quantum state)

- Single copy of  $|\psi\rangle$  given to Anita.

- Her task is to communicate this state to Bharat without destroying the state!

(i.e. without measuring the state!)

- Resource: Anita & Bharat share a Bell state.

$$|\psi\rangle = \alpha_0|0\rangle + \alpha_1|1\rangle \in \mathcal{H}_A,$$

$$|\beta_{00}\rangle = \frac{1}{\sqrt{2}}(|00\rangle + |11\rangle) \in \mathcal{H}_{A_2} \otimes \mathcal{H}_B$$

$$\begin{aligned} \bullet |\psi_0\rangle_{A_1 A_2 B} &= (\alpha|0\rangle + \beta|1\rangle)_{A_1} \otimes \frac{1}{\sqrt{2}}(|00\rangle + |11\rangle)_{A_2 B} \\ &= \frac{1}{\sqrt{2}} [\alpha|0\rangle(|00\rangle + |11\rangle) + \beta|1\rangle(|00\rangle + |11\rangle)] \end{aligned}$$

$$\bullet |\psi_1\rangle_{A_1 A_2 B} = \frac{1}{\sqrt{2}} [\alpha|0\rangle(|00\rangle + |11\rangle) + \beta|1\rangle(|10\rangle + |01\rangle)]$$

$$\begin{aligned} \bullet |\psi_2\rangle_{A_1 A_2 B} &= \frac{1}{2} [\alpha(|0\rangle + |1\rangle)(|00\rangle + |11\rangle) + \beta(|0\rangle - |1\rangle)(|10\rangle + |01\rangle)] \\ &= \frac{1}{2} \left[ |00\rangle_A (\alpha|0\rangle + \beta|1\rangle)_B + |11\rangle_A (\alpha|1\rangle - \beta|0\rangle)_B \right. \\ &\quad \left. + |01\rangle_A (\alpha|1\rangle + \beta|0\rangle)_B + |10\rangle_A (\alpha|0\rangle - \beta|1\rangle)_B \right] \end{aligned}$$

• Anita measures her qubits in  $\{|00\rangle, |11\rangle, |01\rangle, |10\rangle\}$  basis.

|              |           |  |
|--------------|-----------|--|
| $(m_1, m_2)$ | B's state |  |
| $(0, 0)$     |           | $\alpha 0\rangle + \beta 1\rangle =  \psi\rangle!$   |
| $(0, 1)$     |           | $\alpha 1\rangle + \beta 0\rangle = \sigma_x  \psi\rangle = (\sigma_x)^{m_2} (\sigma_z)^{m_1}  \psi\rangle$          |
| $(1, 0)$     |           | $\alpha 0\rangle - \beta 1\rangle = \sigma_z  \psi\rangle = (\sigma_x)^{m_2} (\sigma_z)^{m_1}  \psi\rangle$          |
| $(1, 1)$     |           | $\alpha 1\rangle - \beta 0\rangle = \sigma_x \sigma_z  \psi\rangle = (\sigma_x)^{m_2} (\sigma_z)^{m_1}  \psi\rangle$ |

Note: No faster than light communication!

\* B gets  $|\psi\rangle$  only after A has communicated her measurement outcomes.

\* Without the classical channel, teleportation conveys No information.

History:

Quantum teleportation discovered by Bennett, Brassard, Crépeau, Jozsa, Peres and Wothers, PRL 70, 1895 (1993).

- Experimentally:- Bouwmeester et al (1997) (photon polarization)
- Furusawa et al (1998) ("squeezed" states of light)
- Nielsen et al (1998) (NMR)

(2b) Reduced density operator: Partial trace

Given  $\rho_{AB} \in \mathcal{H}_A \otimes \mathcal{H}_B$ ,

$$\rho_B = \text{Tr}_A(\rho_{AB}) = \sum_i \langle i_A | \rho_{AB} | i_A \rangle,$$

where,  $\{|i_A\rangle\}$  is an ON basis for  $\mathcal{H}_A$ .

$$\text{Similarly } \rho_A = \text{Tr}_B(\rho_{AB}) = \sum_j \langle j_B | \rho_{AB} | j_B \rangle,$$

where,  $\{|j_B\rangle\}$  is an ON basis for  $\mathcal{H}_B$ .

Examples: (Ia)  $\rho_{AB} = |a\rangle\langle a| \otimes |b\rangle\langle b| = |ab\rangle\langle ab|$

$$\rho_A = \langle b|b\rangle |a\rangle\langle a|, \rho_B = \langle a|a\rangle |b\rangle\langle b|$$

(Ib)  $\rho_{AB} = \rho_A \otimes \rho_B$ , where,

$\rho_A$  is a density operator on  $\mathcal{H}_A$

$\rho_B$  is a density operator on  $\mathcal{H}_B$ .

$$\Rightarrow \text{Tr}_A(\rho_{AB}) = \text{Tr}_A(\rho_A) \rho_B = \rho_B$$

$$\text{Tr}_B(\rho_{AB}) = \text{Tr}_B(\rho_B) \rho_A = \rho_A.$$

$$(II) |\Psi_{AB}\rangle = \frac{1}{\sqrt{2}} (|00\rangle + |11\rangle)$$

$$\rho_{AB} = \frac{1}{2} (|00\rangle\langle 00| + |00\rangle\langle 11| + |11\rangle\langle 00| + |11\rangle\langle 11|)$$

$$\begin{aligned} \text{Tr}_A(\rho_{AB}) &= \text{Tr}_B(\rho_{AB}) = \frac{1}{2} (|0\rangle\langle 0| + |1\rangle\langle 1|) \\ &= \frac{I}{2} // \end{aligned}$$

(2c) Teleportation: a local point-of-view!

$$|\psi\rangle_{A_1 A_2 B} = \frac{1}{2} \left[ |00\rangle_A (\alpha|0\rangle + \beta|1\rangle)_B + |11\rangle_A (\alpha|1\rangle - \beta|0\rangle)_B \right. \\ \left. + |10\rangle_A (\alpha|1\rangle + \beta|0\rangle)_B + |11\rangle_A (\alpha|0\rangle - \beta|1\rangle)_B \right]$$

$$\begin{array}{l|l} \alpha|0\rangle + \beta|1\rangle \equiv |\phi_0\rangle_B & \alpha|0\rangle - \beta|1\rangle \equiv |\phi_2\rangle_B \\ \alpha|1\rangle + \beta|0\rangle \equiv |\phi_1\rangle_B & \alpha|1\rangle - \beta|0\rangle \equiv |\phi_3\rangle_B \end{array}$$

$\therefore$  Density operator after Anita's measurement,  
(post-measurement state)

$$\rho_{A_1 A_2 B} = \frac{1}{4} \left[ |00\rangle\langle 00|_A \otimes |\phi_0\rangle\langle \phi_0|_B + |01\rangle\langle 01|_A \otimes |\phi_1\rangle\langle \phi_1|_B \right. \\ \left. + |10\rangle\langle 10|_A \otimes |\phi_2\rangle\langle \phi_2|_B + |11\rangle\langle 11|_A \otimes |\phi_3\rangle\langle \phi_3|_B \right]$$

$\therefore$  Reduced state on Bharat's side:

$$\begin{aligned} \rho_B &= \frac{1}{4} \sum_{i=0}^3 |\phi_i\rangle\langle \phi_i|_B \\ &= \frac{2(|\alpha|^2 + |\beta|^2)}{4} (|0\rangle\langle 0| + |1\rangle\langle 1|) \\ &= \frac{\mathbb{I}}{2}! \end{aligned}$$

$\therefore$  Without classical communication,

$\rho_B$  is completely independent of  $|\psi\rangle$

$\Rightarrow$  B has no information about  $|\psi\rangle$ .